

# Recent Lattice QCD results and implications for BES

Sayantan Sharma



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## Bielefeld-BNL-CCNU collaboration

A. Bazavov, H.-T. Ding, P. Hegde, O. Kaczmarek, F. Karsch, E. Laermann, S. Mukherjee, H. Ohno, P. Petreczky, C. Schmidt, S. Sharma, W. Soeldner, P. Steinbrecher, M. Wagner

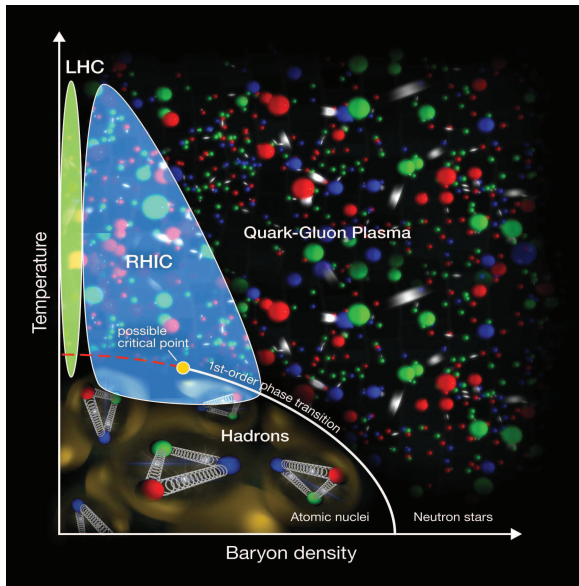
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- 1 The QCD phase diagram: outstanding issues from lattice
- 2 Equation of state at finite  $\mu_B$
- 3 Freezeout and Lattice observables

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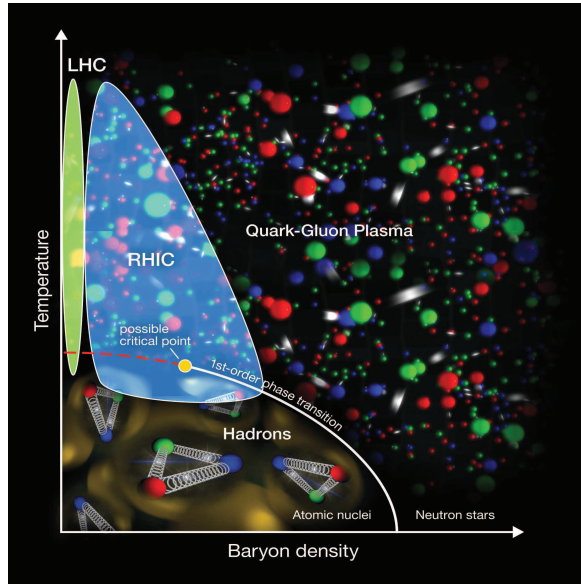
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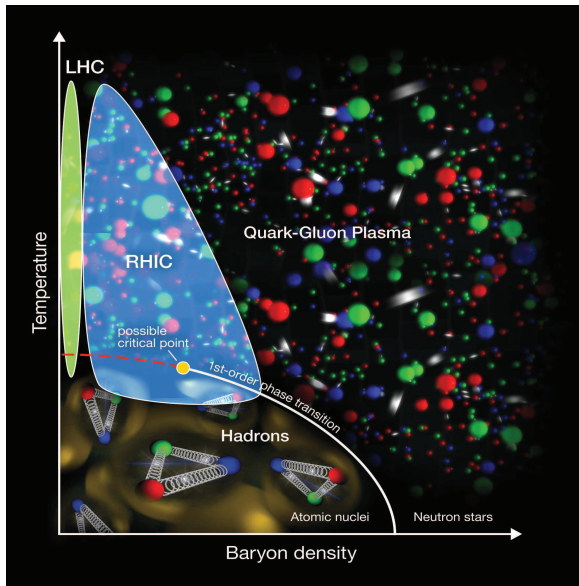
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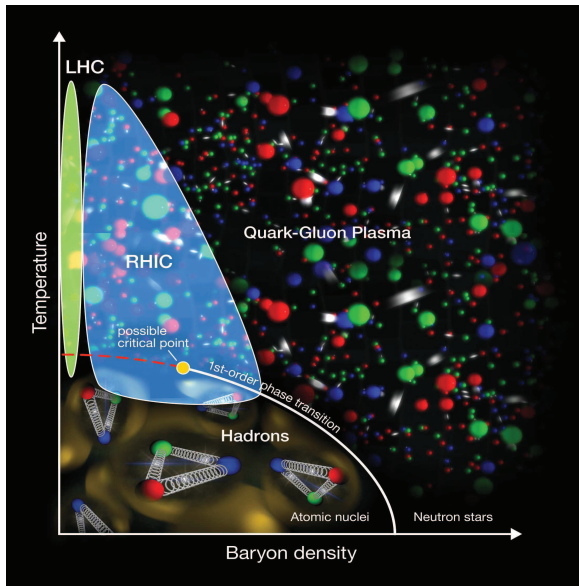
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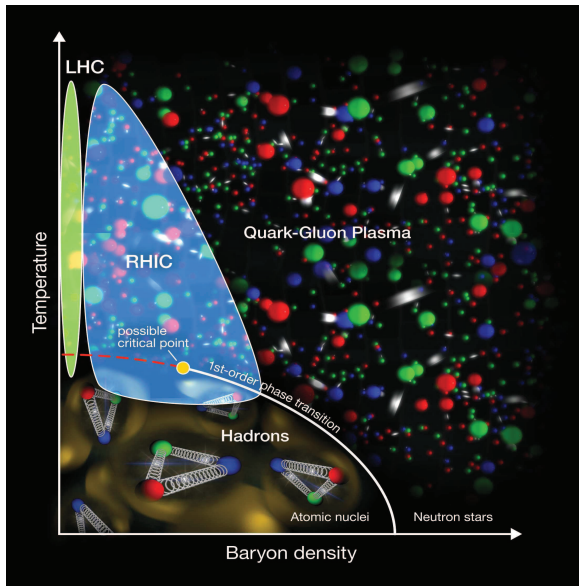
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- Understand what happens to HRG picture at finite  $\mu_B$ .
- Bracket the position of CEP in phase diagram.
- Understand the critical behavior due to the light quarks in the crossover region.





# Basic observables

- One of the methods to circumvent **sign problem** at finite  $\mu$ :  
Taylor expansion of physical observables around  $\mu = 0$  in powers of  $\mu/T$ .

$$\frac{P(\mu_B, T)}{T^4} = \frac{P(0, T)}{T^4} + \frac{1}{2} \left( \frac{\mu_B}{T} \right)^2 \chi_2^B(0, T) + \frac{1}{4!} \left( \frac{\mu_B}{T} \right)^4 \chi_4^B(0) + \dots$$

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- These observables imp. for EoS  $\rightarrow \chi_6^B$  can already constrain QCD pressure in the regime approximated by Hadron Resonance gas model.

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  - Matrix inversions increasing with the order
  - Delicate cancellation between a large number of terms for higher order QNS.

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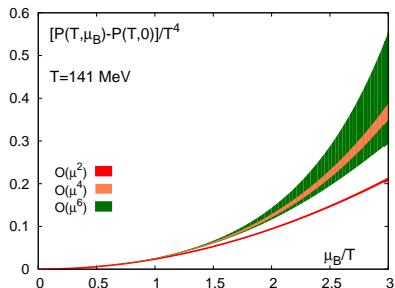
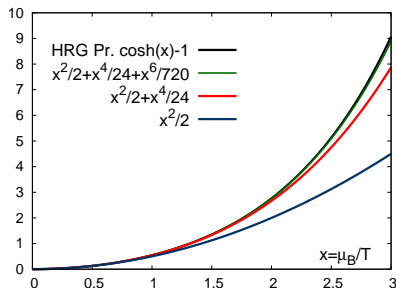
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- Calculating explicitly the lowest eigenvalues improves performance of the fermion inverter. Optimized codes developed to this end.
- Efficient codes based on modern computer architectures are being developed. [O. Kaczmarek, C. Schmidt, P. Steinbrecher, M. Wagner, 14]

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# Constraining EoS

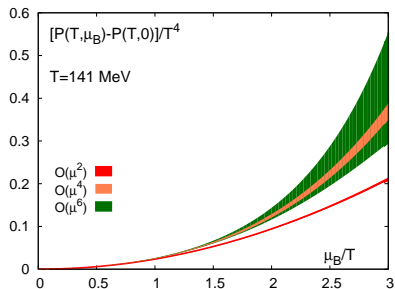
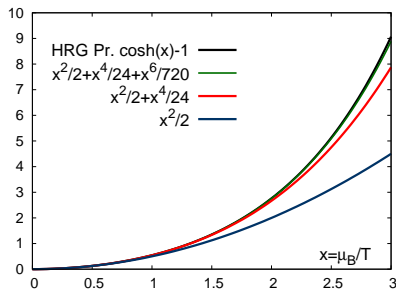
- In a regime where Hadron Resonance gas is anticipated to be a good description of QCD, including  $\chi_6^B$  term already reproduces  $P(\mu_B)$  within 5% accuracy.





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- We are improving the errors on  $\chi_6^B \rightarrow$  increase statistics twofold this year.

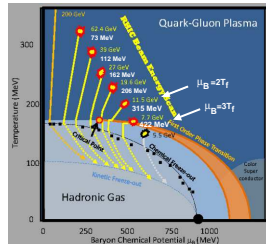
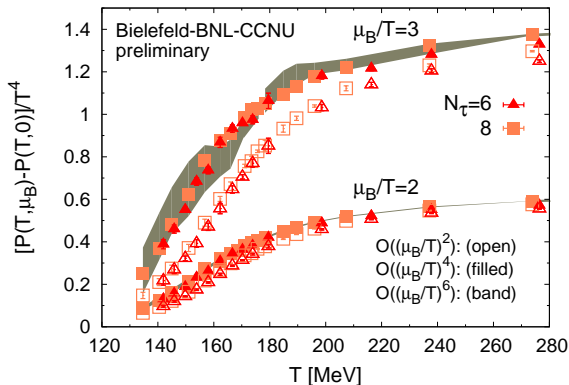


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- Extension to  $\mu_B/T \sim 3$  is in progress.

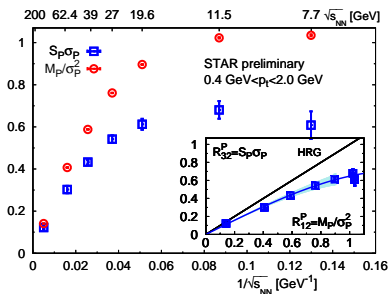


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# The fluctuations near $T_c$

$$\frac{P(T, \mu_B) - P(T, 0)}{T^4} = \frac{\chi_2^B}{2} \left( \frac{\mu_B}{T} \right)^2 \left[ 1 + \frac{1}{12} \frac{\chi_2^B}{\chi_4^B} \left( \frac{\mu_B}{T} \right)^2 + \dots \right].$$

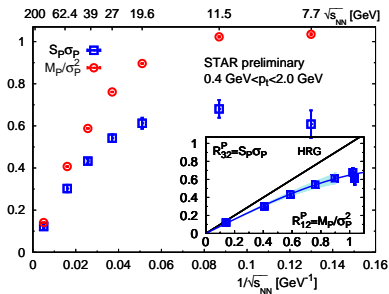


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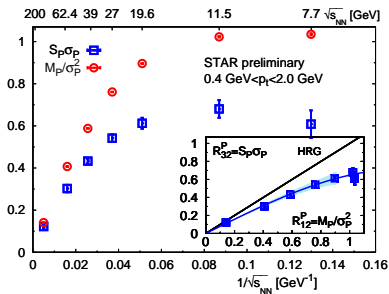
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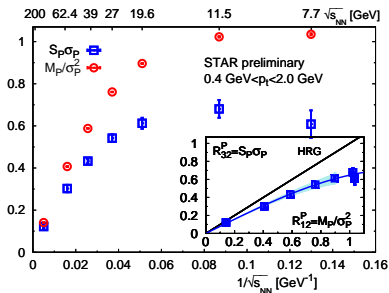
- $S_B \sigma_B = \frac{\mu_B}{T} \frac{\chi_2^B}{\chi_4^B} + \mathcal{O} \left( \frac{\mu_B}{T} \right)^3$   
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- By construction independent of  $\mu_B$ .
- Strangeness neutrality condition:  
 $\frac{n_p}{n_p + n_n} = 0.4.$

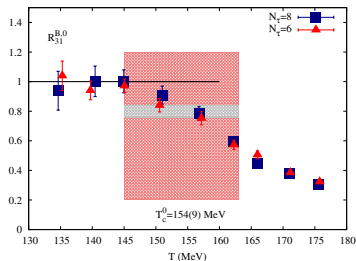
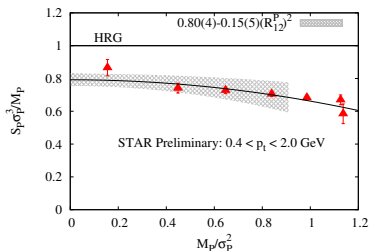


# Is this deviation consistent with the trend from Lattice?

$$\frac{S_B \sigma_B^3}{M_B} = \frac{\chi_2^B}{\chi_4^B} + \frac{1}{6} \left[ \frac{\chi_6^B}{\chi_2^B} - \left( \frac{\chi_2^B}{\chi_4^B} \right)^2 \right] \left( \frac{M_B}{\sigma_B^2} \right)^2.$$

$$R_{31}^B = R_{31}^{B,0} + R_{31}^{B,2} \left( \frac{M_B}{\sigma_B^2} \right)^2. \quad \text{[Karsch et. al., arxiv:1512.06987]}$$

- Experimental data consistent with QCD prediction. **Caveat**  $n_P \neq n_B$ !

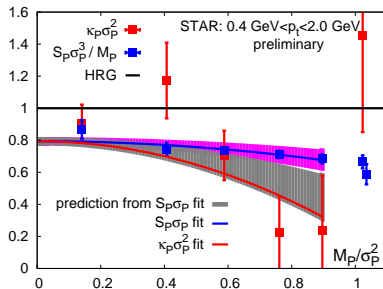


# More Observables?

$$\text{NLO } \kappa_B \sigma_B^2 \text{ for } \mu_Q \sim \mu_S \sim 0: R_{42}^{B,2} = \frac{1}{2} \left[ \frac{\chi_6^B}{\chi_2^B} - \left( \frac{\chi_2^B}{\chi_4^B} \right)^2 \right] = 3R_{31}^{B,2}.$$

- Fit to experimental data shows these quantities are closely related.
- $R_{31}^{B,0} \approx R_{42}^{B,0}$ . At NLO consistent within large errors in the data. [

Bielefeld-BNL-CCNU collaboration, In preparation, 16]

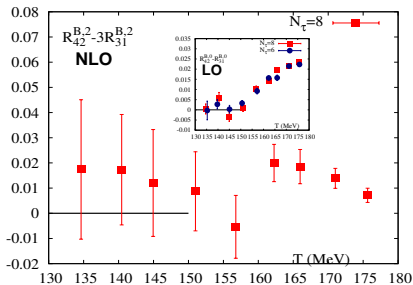


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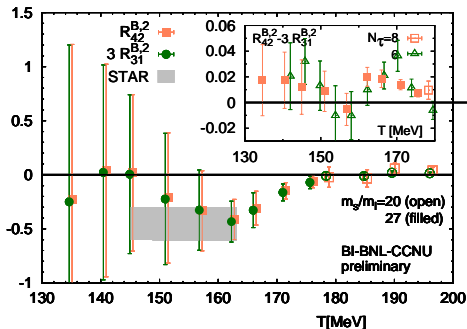
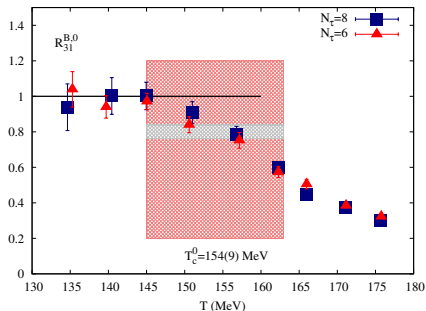
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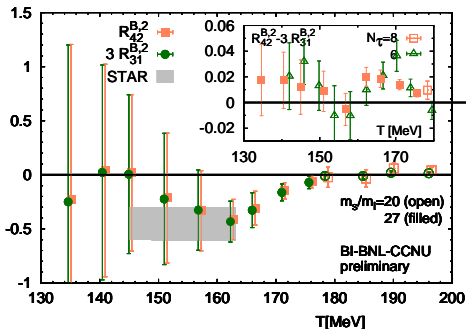
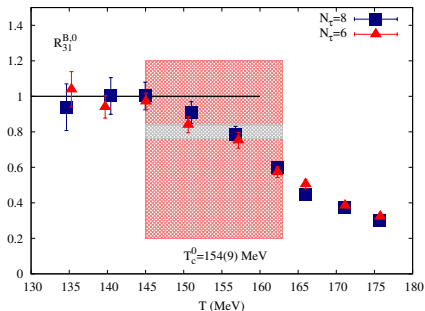
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- Aiming for a factor two reduction of errors on 6th order cumulants near  $T_c$ .



# Freezeout curve: Input from lattice

- Freezeout curve parametrized as  $T = T_{f,0}(1 - \kappa_2^f \mu_B^2 / T_{f,0}^2)$ .

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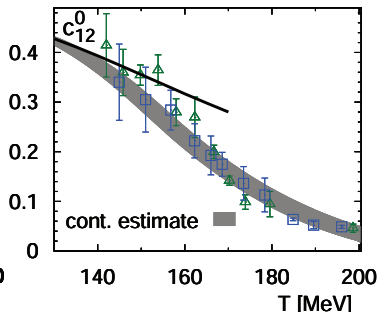
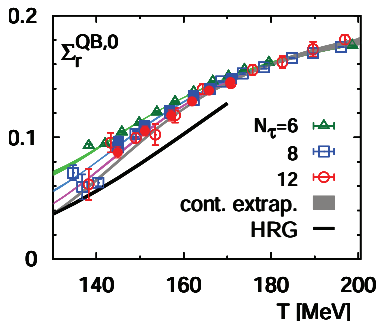
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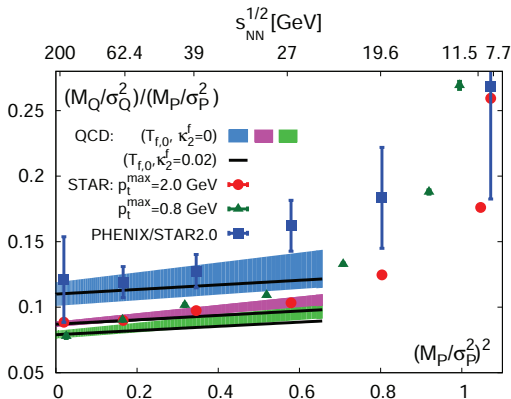
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- As a result  $\Sigma_r^{QB}(\mu_B) = \Sigma_r^{QB}(0) \left[ 1 + c_{12} (R_{12}^B)^2 \right] + \mathcal{O} (R_{12}^B)^4$ .

# Freezeout curve: Input from lattice

- Freezeout curve parametrized as  $T = T_{f,0}(1 - \kappa_2^f \mu_B^2 / T_{f,0}^2)$ .
- Basic observables  $\Sigma_r^{QB} = \frac{R_{12}^Q}{R_{12}^B}$  ,  $R_{12}^X = \frac{\chi_1^X}{\chi_2^X}$ .
- Expanding the observable about the freezeout surface at  $\mu_B = 0$ ,  
 $\Sigma_r^{QB}(\mu_B) = \Sigma_r^{QB}(0) + \left[ \Sigma_r^{QB,2} - \kappa_2^f T_{f,0} \frac{d\Sigma_r^{QB,0}}{dT} \Big|_{T_{f,0}} \right] \frac{\mu_B^2}{T^2}$ .
- Instead from experimental parametrization, get  $\mu_B$  from the first order Taylor expansion of  $R_{12}^B$  as  $\mu_B = T \frac{R_{12}^{B,0}}{R_{12}^{B,1}}$ .
- As a result  $\Sigma_r^{QB}(\mu_B) = \Sigma_r^{QB}(0) \left[ 1 + c_{12} (R_{12}^B)^2 \right] + \mathcal{O} (R_{12}^B)^4$ .
- An estimate of  $\Sigma_r^{QB}$  and  $R_{12}^B$  from experiments allows us to calculate  $c_{12}$ . [ Bielefeld-BNL-CCNU collaboration, 15]

- **Caveat:** In experiments one measures protons  $\Sigma_r^{QP}$ ,  $R_{12}^P$ . Need to understand proton vs baryon number distributions. [Asakawa & Kitazawa, 12]. Within HRG at least  $R_{12}^B$  is mimicked by  $R_{12}^P$  within 10%.
- Additionally take into account also corrections due to finite range of momenta of detected particles.  
[Karsch, Morita and Redlich, 15, P Garg et. al., 13, Bzdak & Koch, 12].
- From the 2 independent expressions of  $\Sigma_r^{QB}$  we extract  $c_{12}(T_{f,0}, \kappa_2^f) = c_{12}(T_{f,0}) - \kappa_2^f D_{12}$ .





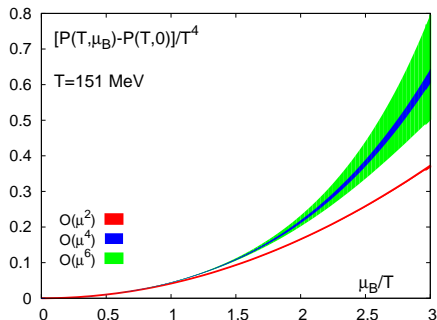
- This exercise give  $T_{f,0} = 147$  MeV consistent with expectation that its at or below  $T_c$ .

Curvature:  $\kappa_2^f < -0.012(15) \rightarrow$  near to chiral curvature  $\kappa_2^B = 0.0066(7)$ .

[ Bielefeld-BNL-CCNU collaboration, 15]

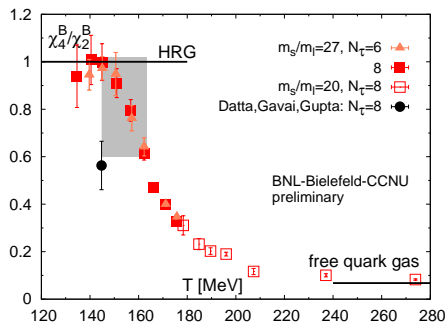
# Breakdown of HRG

- Breakdown of HRG+ onset of criticality can be already constrained with  $\chi_6^B$ .
- Near critical point all terms in the Taylor expansion nearly equal  $\rightarrow$  need to improve the errors to observe!
- At CEP:  $\chi_n > 0$ ,  $\kappa_B \sigma_B^2 > 1$



# Critical-end point search from Lattice

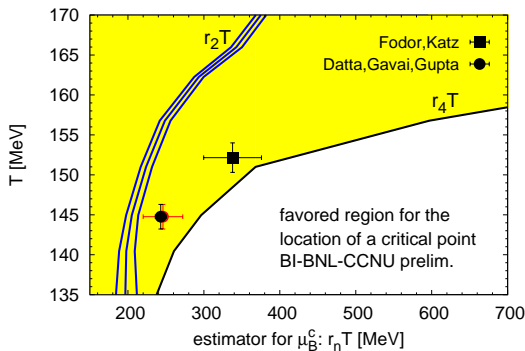
- Radius of convergence:  $r_{2n} \equiv \sqrt{2n(2n-1)} \left| \frac{\chi_{2n}^B}{\chi_{2n+2}^B} \right|$ .
- Only existing result  $T_{CEP} = 0.94 T_c$   $\mu/T = 1.68(5)$   
[ S.Datta, R. Gavai, S.Gupta, 13, Mumbai group]
- Lowest  $r_{n=2}$  varies significantly from our estimates and HRG  $\rightarrow$  lattice cut-off effects needs to be considered!



# Critical-end point search from Lattice

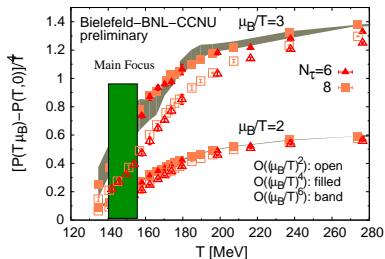
- Current errors on our  $\chi_6^B/\chi_4^B$  only allow us to define a favored region for CEP!
- $\chi_8^B$  measured to get errors bounds on radius of convergence estimates.
- Connection to experiments non-trivial due to non-equilibrium effects.

[ S. Mukherjee, Y. Yin, R. Venugopalan, 15]



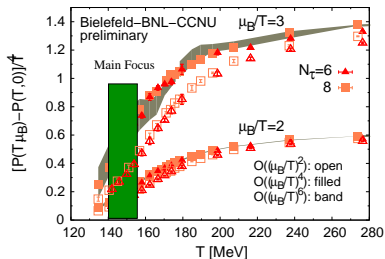


# Outlook



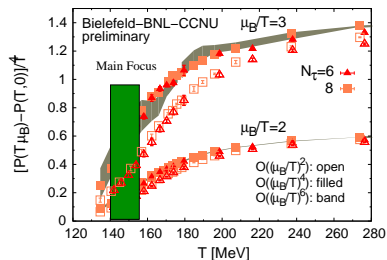
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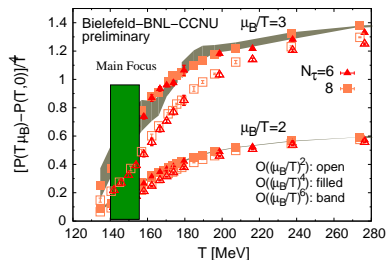
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- Analysis of  $\chi_8^B$  ongoing. Crucial to estimate the errors on the EoS measured with the sixth order cumulants.
- Higher order cumulants will also help in bracketing the possible CEP. Current LQCD data suggest it is  $\mu_B/T \leq 2$  but a larger value cannot be ruled out.